

Inertia-gravity waves in geophysical vortices

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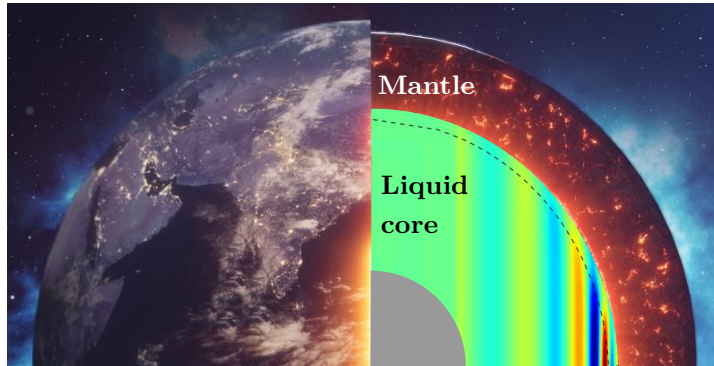
European
Research
Council

INTRODUCTION

An interdisciplinary collaboration

Myself = **Fluid dynamicist**

Y. Colin de Verdière



+ **Astrophysical** flows

+ **Lab. experiments**

Series of papers (submitted)

Colin de Verdière & **Vidal**, 2023, *The spectrum of the Poincaré operator in an ellipsoid*, [arXiv:2305.01369](https://arxiv.org/abs/2305.01369)

Vidal & Colin de Verdière, 2023, *Inertia-gravity waves in geophysical vortices*

INTRODUCTION

Vortices in geophysical fluid dynamics

$$\text{Vortex} \Rightarrow \text{Vorticity } \boldsymbol{\omega} = \nabla \times \boldsymbol{v} \neq \mathbf{0}$$

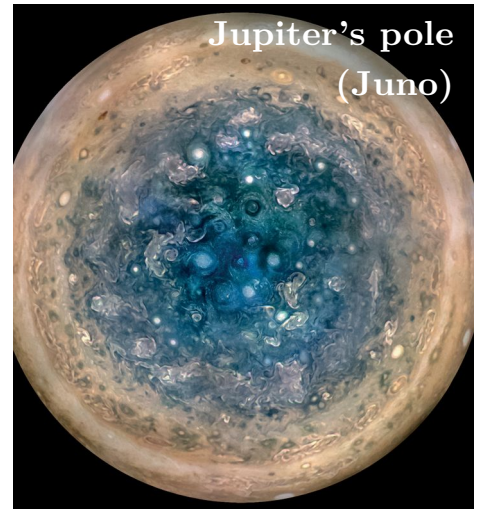
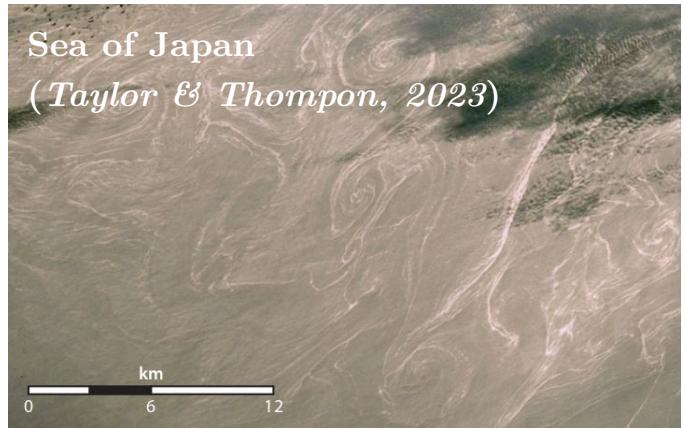
Definition: “What is [a vortex] ? It is like pornography. It is hard to defined but if you see it, you recongnise it immediately.”

G. K. Vallis

< m



~ km



At all length scales!

INTRODUCTION

Standard equations for geophysical flows

$$\partial_t \mathbf{v} + 2\boldsymbol{\Omega} \times \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v} + d_t \boldsymbol{\Omega} \times \mathbf{r} = -\nabla p + (\rho/\rho_*) \mathbf{g} + \nu \nabla^2 \mathbf{v} + \mathbf{f}$$

$$\nabla \cdot \mathbf{v} = 0$$

$$\partial_t \rho + \mathbf{v} \cdot \nabla \rho_0 = \kappa \nabla^2 \rho$$

+ Boundary Conditions (BC)

Background density

Boussinesq equations

2 key ingredients!

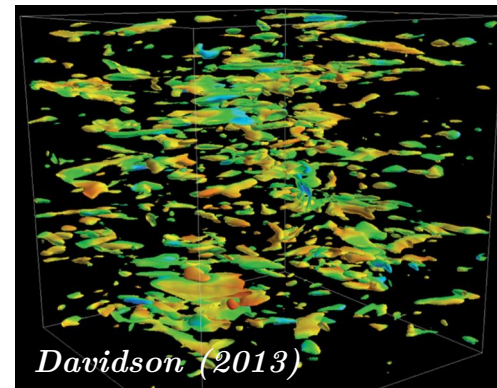
- Global rotation
- Density stratification

$R \gg S$



Elongated vortices

$R \ll S$



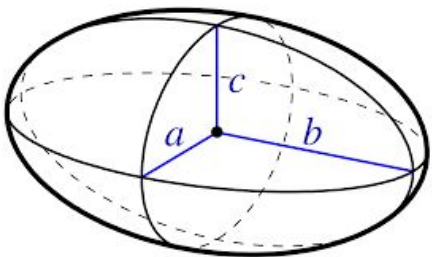
Pancake-like vortices

OR

INTRODUCTION

Dynamics of ellipsoidal vortices

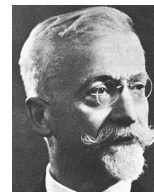
- (Nearly) **isolated** vortices
- Fluid confined within an **ellipsoid**
- **Linear** wave motions



1880 - 1930



Poincaré

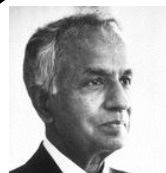


Cartan

1800 - 1880

Dirichlet, Jacobi, Riemann...

1960 - 1990



Chandrasekhar
(Nobel prize)



Lebovitz

Friedlander



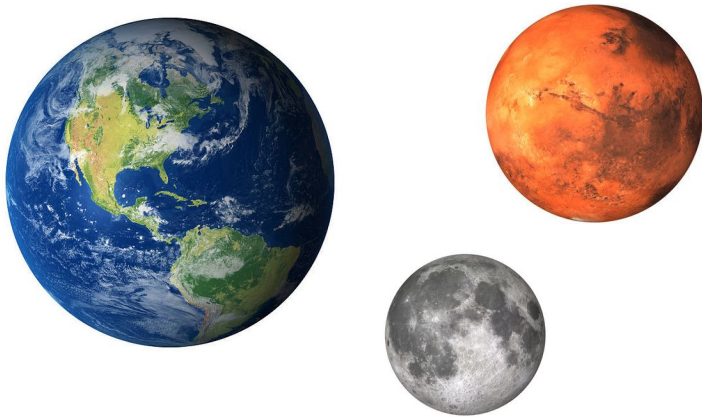
Long-standing history!

OUTLINE

2 **wave** families, 2 different **motivations** (but **similar** maths)

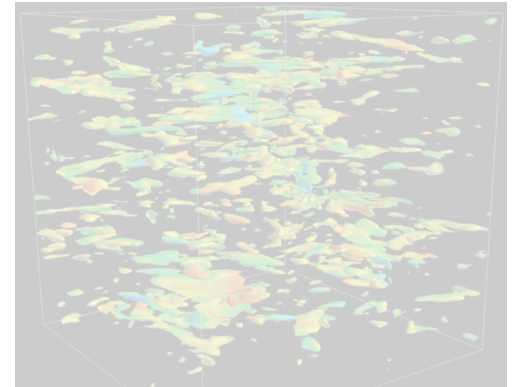
Paper I (CdV&Vidal)

Inertial waves



Paper II (Vidal&CdV)

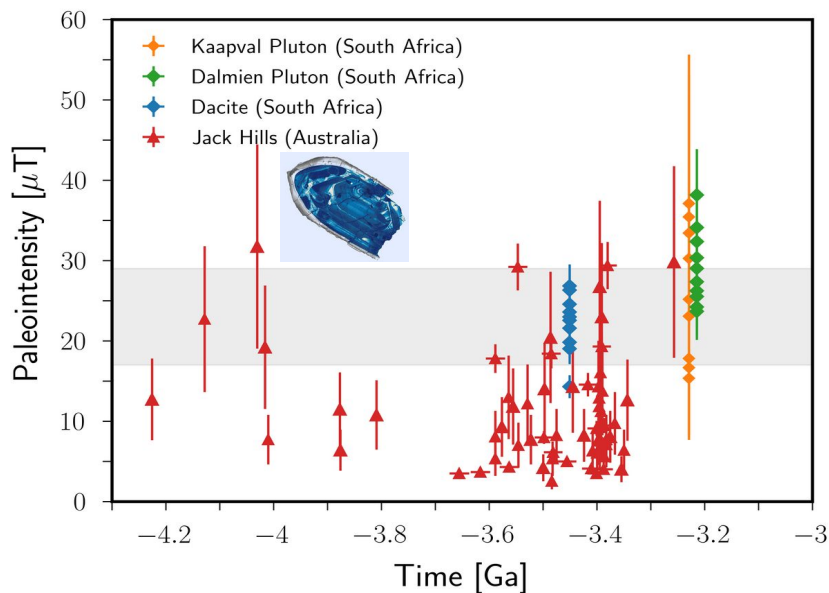
Inertia-gravity waves



INERTIAL WAVES

Geophysical motivation (1/2)

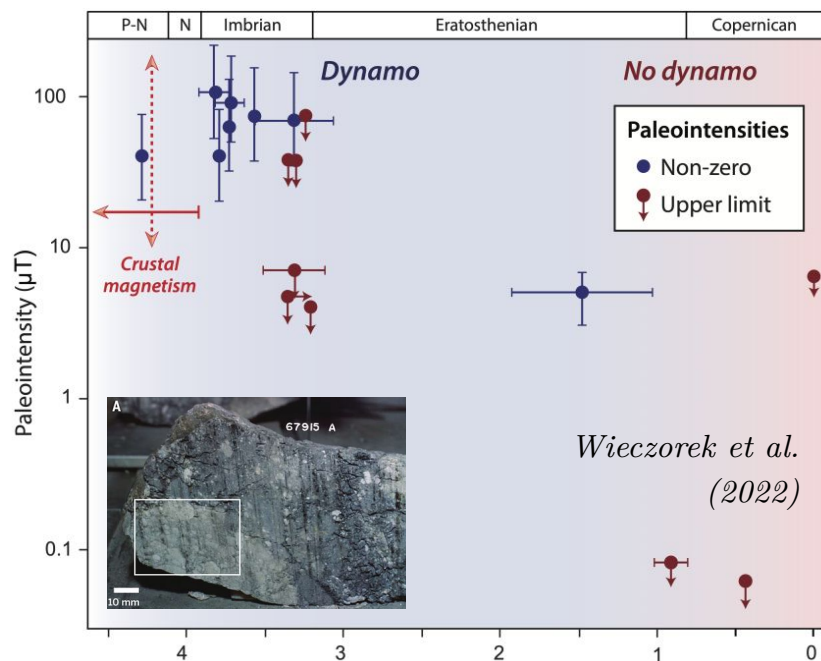
Earth



Tarduno et al. (2015, 2020, 2023)

Borlina et al. (2020)

Moon



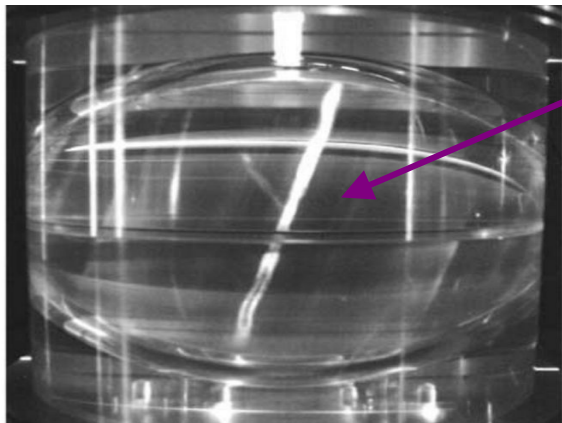
Origin of these dynamos: wave turbulence?

INERTIAL WAVES

Geophysical motivation (2/2)

Orbitally driven (forced) flow

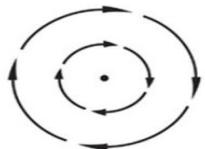
Noir et al. (2003)



“Poincaré” flow

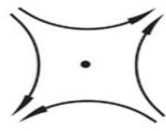
Rotation

Strain



Sphere

+



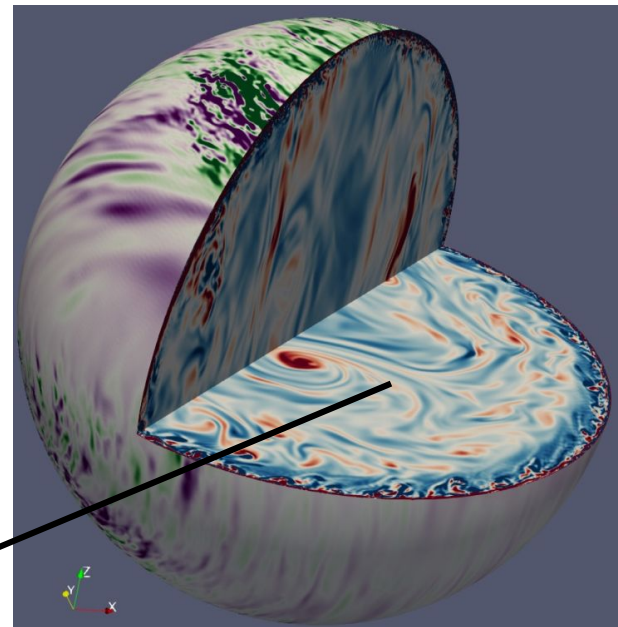
Ellipsoid

Nonlinear interactions

$$(\mathbf{v} \cdot \nabla)\mathbf{v}$$

Turbulence

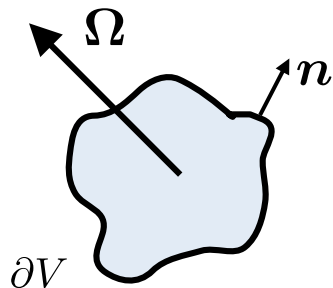
Cébron et al. (2019)



Wave turbulence

INERTIAL WAVES

Spectral problem (1/2)

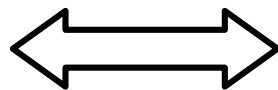


$$[\mathbf{v}, p] = [\mathbf{u}, \Phi] \exp(i\omega t)$$

$$i\omega \mathbf{u} + 2\boldsymbol{\Omega} \times \mathbf{u} = -\nabla \Phi$$

$$\nabla \cdot \mathbf{u} = 0$$

$$\mathbf{u} \cdot \mathbf{n}|_{\partial V} = 0$$



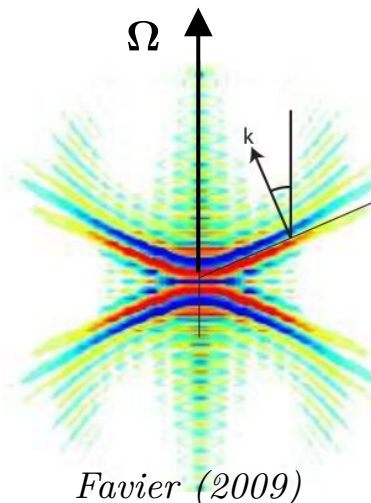
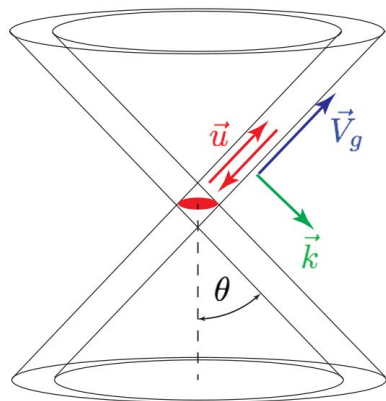
$$\omega^2 \nabla^2 \Phi = (2\boldsymbol{\Omega} \cdot \nabla)^2 \Phi.$$

$$\omega^2 \nabla \Phi \cdot \mathbf{n} = i\omega (2\boldsymbol{\Omega} \times \mathbf{n}) \cdot \nabla \Phi + 4(\boldsymbol{\Omega} \cdot \nabla \Phi)(\boldsymbol{\Omega} \cdot \mathbf{n}) \text{ on } \partial V$$

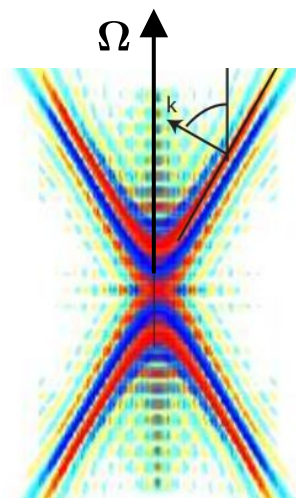
Principal symbol

$$p := \omega^2 |\mathbf{k}|^2 - (2\boldsymbol{\Omega} \cdot \mathbf{k})^2$$

$$\begin{cases} |\omega| < 2|\boldsymbol{\Omega}| : & \text{Hyperbolic} \\ |\omega| > 2|\boldsymbol{\Omega}| : & \text{Elliptic} \end{cases}$$

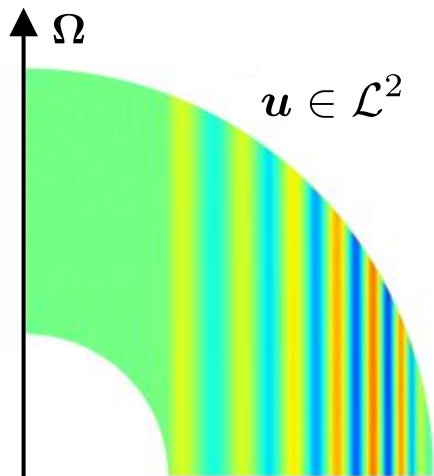


Favier (2009)



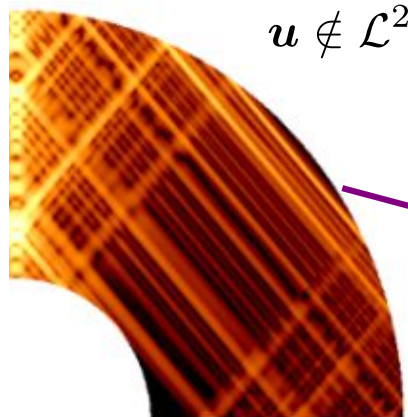
INERTIAL WAVES

Spectral problem (2/2)



$u \in \mathcal{L}^2$

Vidal (2018)



$u \notin \mathcal{L}^2$

Rieutord et al. (2001)

An ill-posed (self-adjoint) problem

- (Almost empty) **point** spectrum
- **Continuous** spectrum

Attractors for Two-Dimensional Waves with Homogeneous Hamiltonians of Degree 0

YVES COLIN DE VERDIÈRE
Université Grenoble-Alpes, Institut Fourier

LAURE SAINT-RAYMOND
Ecole Normale Supérieure de Lyon, UMPA

Abstract

In domains with topography, inertial and internal waves exhibit interesting features. In particular, numerical and lab experiments show that, in two dimensions, for generic forcing frequencies, these waves concentrate on attractors. The goal of this paper is to analyze mathematically this behavior, using tools from spectral theory and microlocal analysis. © 2019 Wiley Periodicals, Inc.

! Shell \neq Ellipsoid !

INERTIAL WAVES

Polynomial description in an ellipsoid

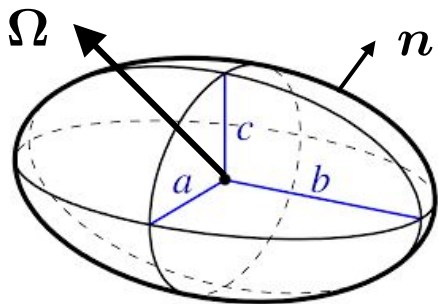
\mathcal{V} : Hilbert space of square-integrable vector fields

$$\langle \mathbf{u}_1, \mathbf{u}_2 \rangle := \int_V \mathbf{u}_1^\dagger \cdot \mathbf{u}_2 \, dV$$

$$\mathcal{V}^0 := \{ \mathbf{u} \in \mathcal{V}, \nabla \cdot \mathbf{u} = 0 \text{ in } V, \mathbf{u} \cdot \mathbf{n} |_{\partial V} = 0 \}$$

\mathcal{P}_n : Vector polynomial functions whose components $\propto x^i y^j z^k$ with $i + j + k \leq n$

$$\mathcal{V}_n^0 := \mathcal{V}^0 \cap \mathcal{P}_n, \quad \dim(\mathcal{V}_n^0) = n(n+1)(2n+7)/6$$



Poincaré operator (bounded & self-adjoint)

$$i\mathcal{C}(\mathbf{u}) := i\mathbb{L}(2\boldsymbol{\Omega} \times \mathbf{u}), \quad \mathbb{L} : \mathcal{V} \rightarrow \mathcal{V}^0$$

$$(\partial V) : \left(\frac{x}{a}\right)^2 + \left(\frac{y}{b}\right)^2 + \left(\frac{z}{c}\right)^2 = 1$$

$$\mathcal{C}|_{\mathcal{V}_n^0} \subseteq \mathcal{V}_n^0$$

$$\bigoplus_n \mathcal{V}_n^0 \text{ dense in } \mathcal{V}^0$$

Pure point spectrum in ellipsoids

INERTIAL WAVES

Microlocal analysis

Asymptotic measure

$$\int_{-\infty}^{\lambda} d\pi_{\infty} = \frac{1}{4\pi} \text{Area}(\mathfrak{S}_{\lambda} \cap S^2)$$

$$\mathfrak{S}_{\lambda} := \{\tilde{\mathbf{k}} \in \mathbb{R}^3 \mid \omega(\tilde{\mathbf{k}}) \leq \lambda\}$$

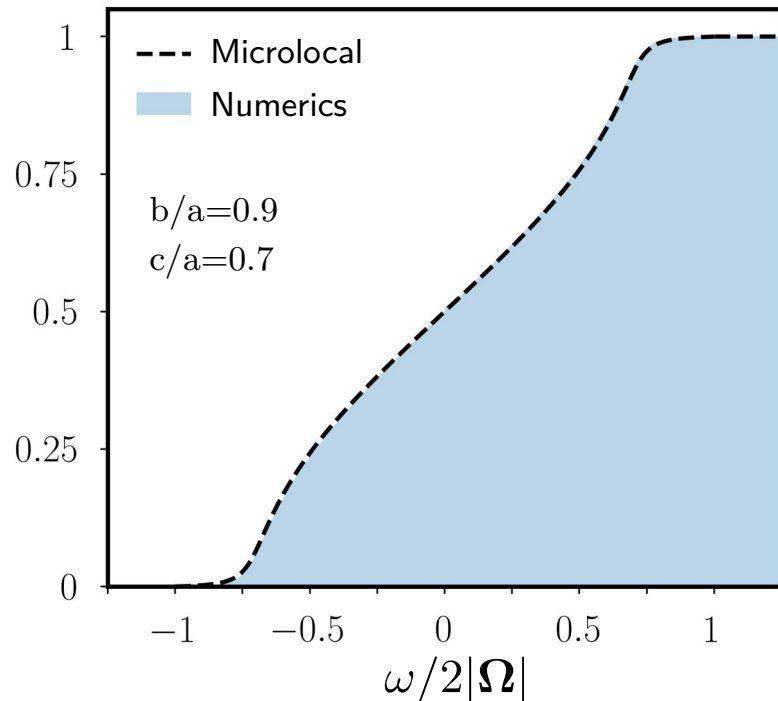
$$\tilde{\mathbf{k}} = (k_x/a, k_y/b, k_z/c)^{\top}$$

Rescaled **dispersion relation** of inertial waves

Proof:

- Pseudo-differential operator of degree 0
- Boutet de Monvel's calculus

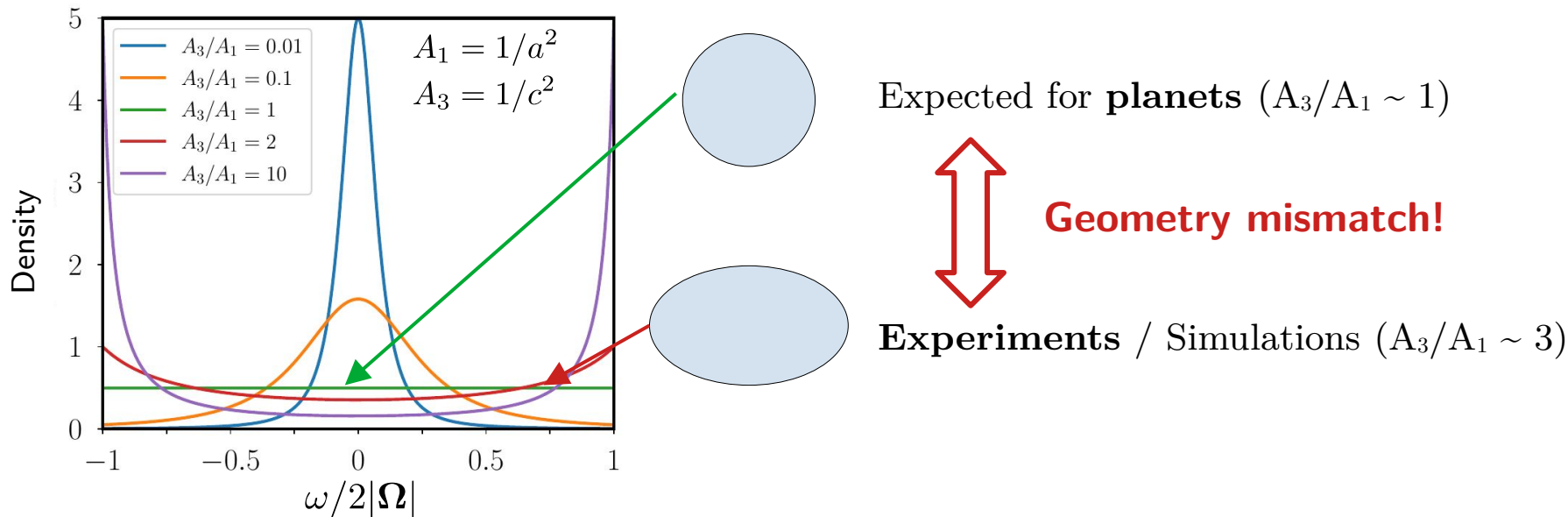
...



Agreement with numerics

INERTIAL WAVES

Physical implications & mathematical perspectives



Beyond ellipsoids?

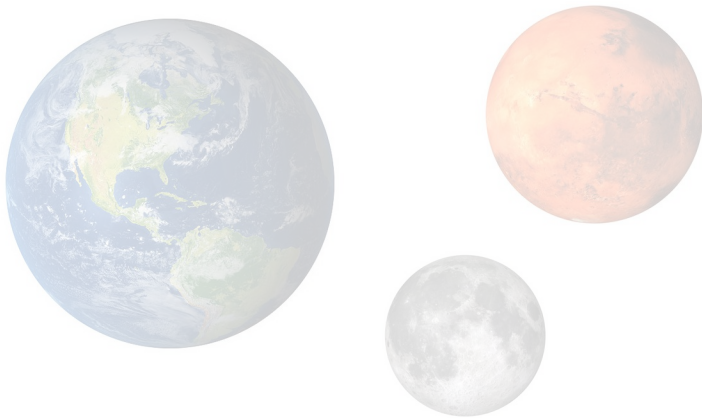
- **Continuous** spectrum: attractors' properties?
- (Almost empty) **point** spectrum: **perturbation** theory?

OUTLINE

2 **wave** families, 2 different **motivations** (but **similar** maths)

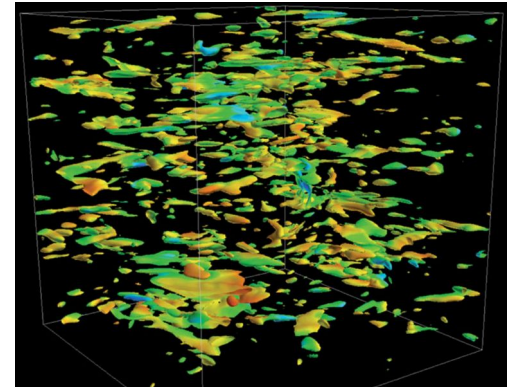
Paper I (CdV&Vidal)

Inertial waves



Paper II (Vidal&CdV)

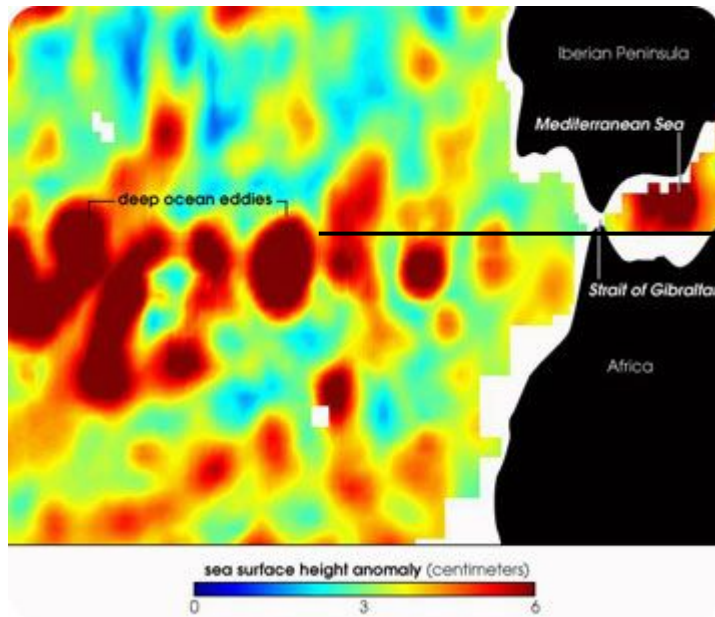
Inertia-gravity waves



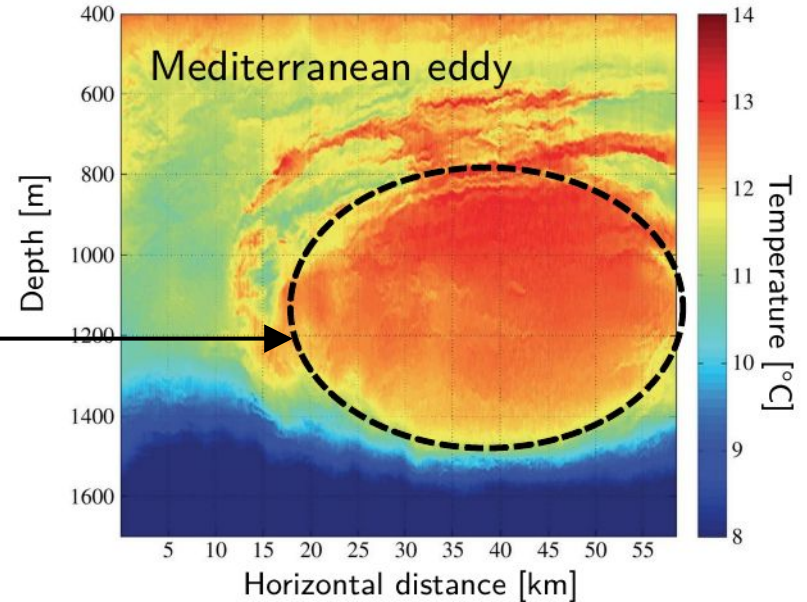
INERTIA-GRAVITY WAVES

Geophysical vortices (1/2)

“Meddies”



Yan et al. (2006)



- **Anticyclonic** vortices ($\rightarrow 50 \text{ cm.s}^{-1}$)
- 10^9 - 10^{11} tons of **salt** per eddy

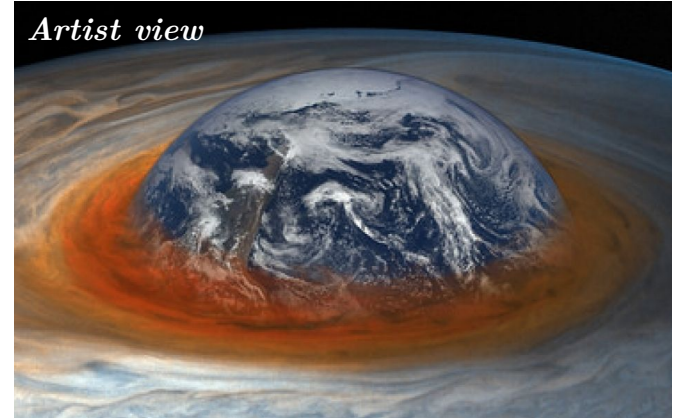
INERTIA-GRAVITY WAVES

Geophysical vortices (2/2)

Jovian vortices (e.g. GRS)



Juno - NASA - JPL-Caltech

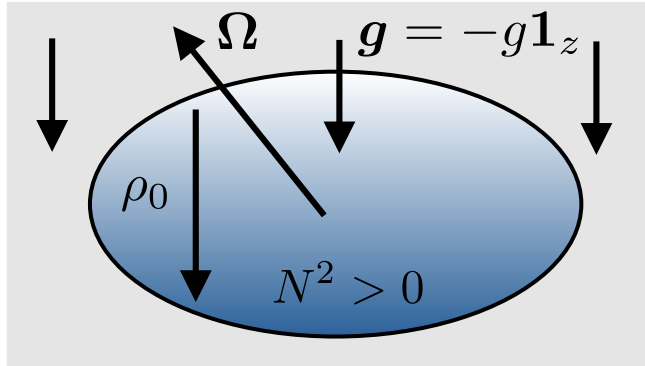


- Horizontal: $10^3 - 10^4$ km
- Vertical: a few hundred of km

Strongly flattened!

INERTIA-GRAVITY WAVES

Nonlinear spectral problem in ellipsoids (1/3)



$$N^2 = \mathbf{g} \cdot \nabla \rho_0 / \rho_*$$

Brunt-Väisälä
frequency

- $N^2 > 0$: Stably stratified (i.e. light above dense)
- $N^2 < 0$: Unstable

Primitive equations

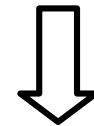
$$\partial_t \mathbf{v} + 2\boldsymbol{\Omega} \times \mathbf{v} = -\nabla p + (\rho/\rho_*) \mathbf{g}$$

$$\nabla \cdot \mathbf{v} = 0$$

$$\partial_t \rho / \rho_* = (\mathbf{v} \cdot \mathbf{1}_z) N^2 / g$$

Wave-like equation

$$\partial_{tt}^2 \mathbf{v} + 2\boldsymbol{\Omega} \times (\partial_t \mathbf{v}) + N^2 (\mathbf{v} \cdot \mathbf{1}_z) \mathbf{1}_z = -\nabla \partial_t p$$

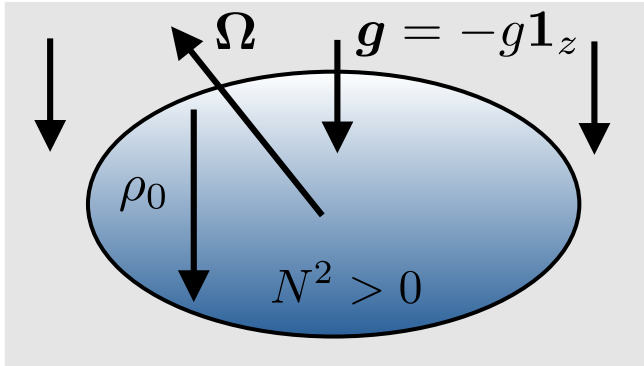


$$[\mathbf{v}, p] = [\mathbf{u}, \Phi] \exp(i\omega t)$$

$$-\omega^2 \mathbf{u} + 2\boldsymbol{\Omega} \times (i\omega \mathbf{u}) + N^2 (\mathbf{u} \cdot \mathbf{1}_z) \mathbf{1}_z = -i\omega \nabla \Phi$$

INERTIA-GRAVITY WAVES

Nonlinear spectral problem in ellipsoids (2/3)



$$-\omega^2 \mathbf{u} + \underbrace{\omega i \mathbb{L}(2\boldsymbol{\Omega} \times \mathbf{u})}_{i\mathcal{C}(\mathbf{u})} + \underbrace{\mathbb{L}(N^2(\mathbf{u} \cdot \mathbf{1}_z)\mathbf{1}_z)}_{\mathcal{K}(\mathbf{u})} = \mathbf{0}$$

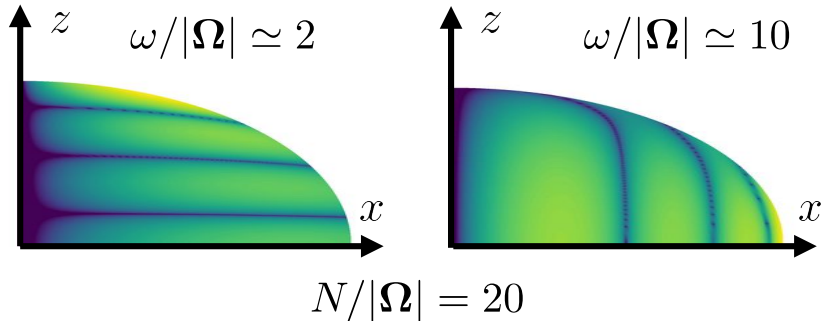
$$\mathcal{C}|_{\mathcal{V}_n^0} \subseteq \mathcal{V}_n^0$$

$$\mathcal{K}|_{\mathcal{V}_n^0} \subseteq \mathcal{V}_n^0$$

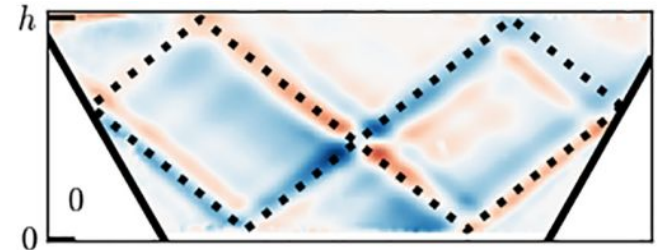
$$\bigoplus_n \mathcal{V}_n^0 \text{ dense in } \mathcal{V}^0$$

Pure point spectrum in ellipsoids

Square-integrable



Attractors (e.g. in lab. experiment)



Pacary et al. (2023)

INERTIA-GRAVITY WAVES

Nonlinear spectral problem in ellipsoids (2/2)

Principal symbol

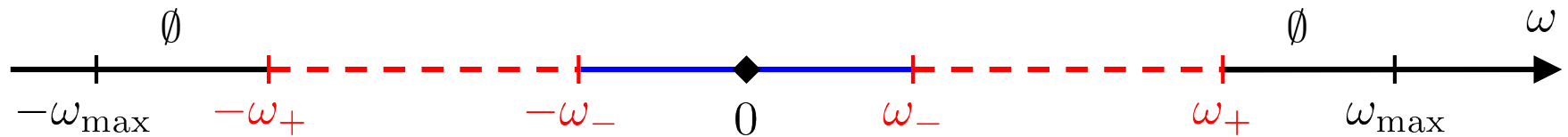
$$\mathfrak{p} := |\mathbf{k}|^2 \omega^2 - [N^2 |\mathbf{1}_z \times \mathbf{k}|^2 + (2\boldsymbol{\Omega} \cdot \mathbf{k})^2]$$

$$\begin{cases} \omega_-^2 < \omega^2 < \omega_+^2 : & \text{Hyperbolic} \\ 0 < \omega^2 < \omega_-^2 : & \text{Elliptic} \end{cases}$$

$$2\omega_{\pm}^2 = [N^2 + 4|\boldsymbol{\Omega}|^2] \pm \sqrt{[N^2 + 4|\boldsymbol{\Omega}|^2]^2 - 16N^2(\boldsymbol{\Omega} \cdot \mathbf{1}_z)^2}$$

— Elliptic in V
- - - Hyperbolic in V

2 wave families?



INERTIA-GRAVITY WAVES

Hyperbolic interval (1/2)

Asymptotic measure ($\omega_- < |\omega| < \omega_+$)

$$\int_{-\infty}^{\lambda} d\pi_{\infty} = \frac{1}{8\pi} \text{Area}(\mathfrak{S}_{\lambda} \cap S^2)$$

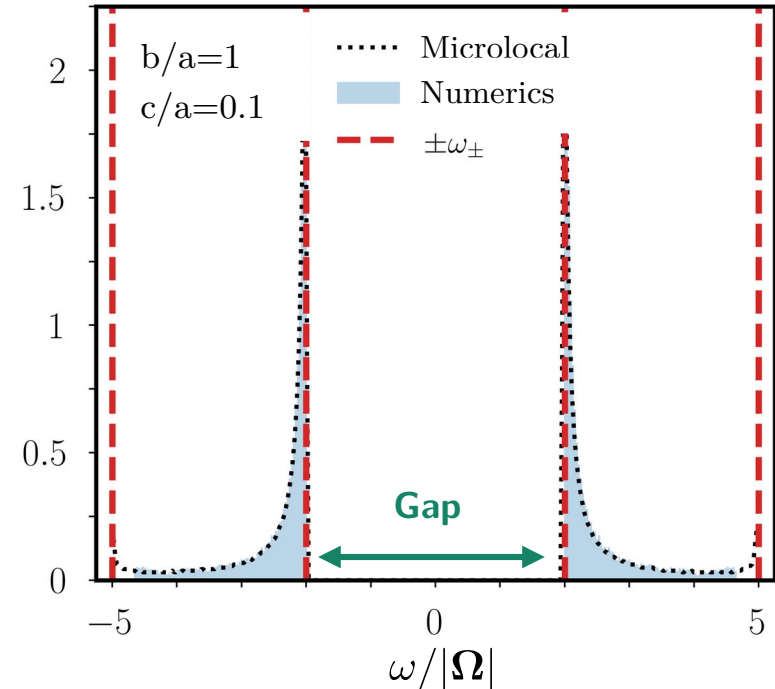
$$\mathfrak{S}_{\lambda} := \{\tilde{\mathbf{k}} \in \mathbb{R}^3 \mid \omega(\tilde{\mathbf{k}}) \leq \lambda\}$$

$$\tilde{\mathbf{k}} = (k_x/a, k_y/b, k_z/c)^{\top}$$

Rescaled **dispersion relation** of inertia-gravity waves

Proof: - Weyl asymptotics of the equivalent matrix operator

$$\mathcal{L} := \begin{pmatrix} 0 & \mathcal{I} \\ \mathcal{K} & i\mathcal{C} \end{pmatrix}$$

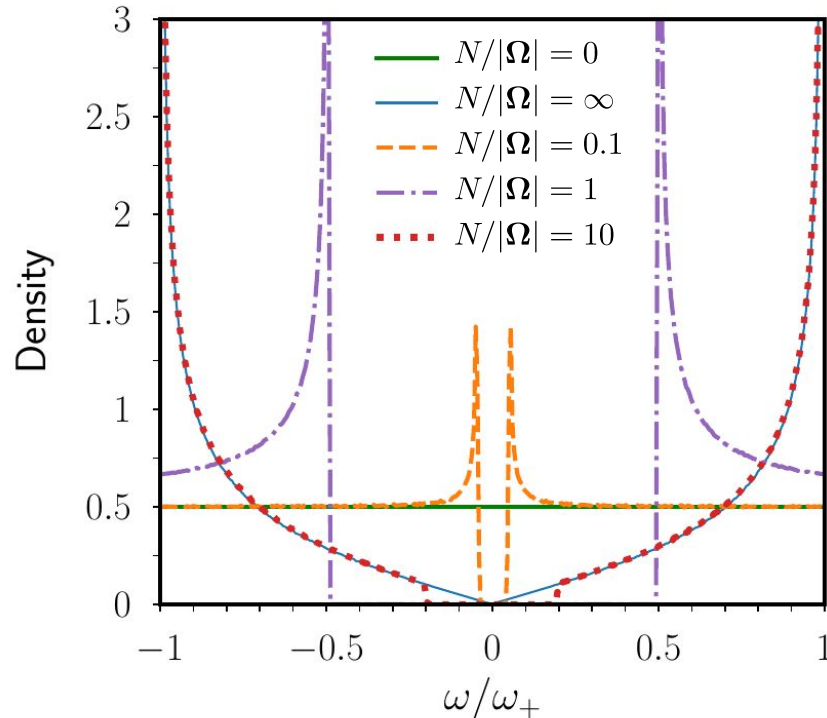


Agreement with numerics

INERTIA-GRAVITY WAVES

Hyperbolic interval (2/2)

Ball $b/a=c/a=1$



Some differences with pure inertial waves

- **Non-uniform** distribution in the ball
- Gap size function of the orientation of $\mathbf{\Omega}$

$$2\omega_{\pm}^2 = [N^2 + 4|\mathbf{\Omega}|^2] \pm \sqrt{[N^2 + 4|\mathbf{\Omega}|^2]^2 - 16N^2(\mathbf{\Omega} \cdot \mathbf{1}_z)^2}$$

- Different **orthogonality** conditions for $(\mathbf{u}_1, \mathbf{u}_2)$

$$\omega_1 \neq \omega_2$$

IW

$$\langle \mathbf{u}_2, \mathbf{u}_1 \rangle = 0$$

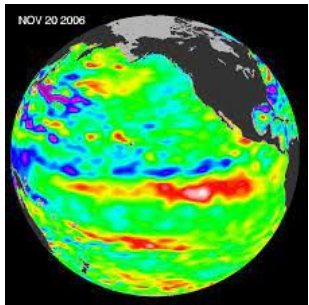
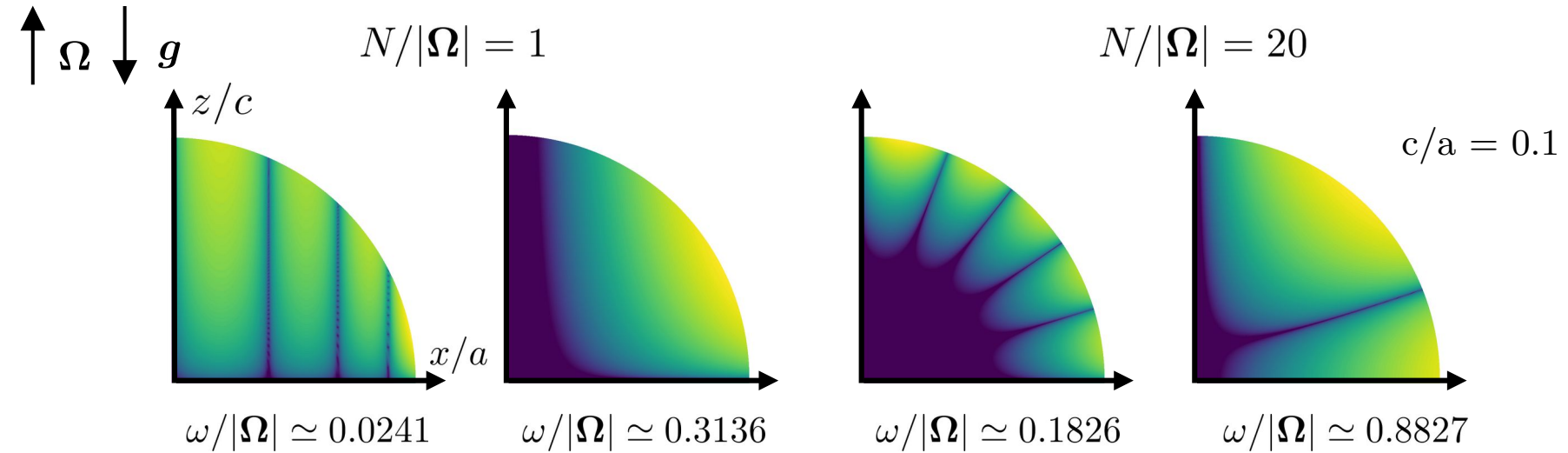
IGW

$$(\omega_2 + \omega_1) \langle \mathbf{u}_2, \mathbf{u}_1 \rangle = \langle \mathbf{u}_2, i\mathcal{C}(\mathbf{u}_1) \rangle$$

$$\omega_1 \omega_2 \langle \mathbf{u}_2, \mathbf{u}_1 \rangle = -\langle \mathbf{u}_2, \mathcal{K}(\mathbf{u}_1) \rangle$$

INERTIA-GRAVITY WAVES

Elliptic interval (1/2)



- **Similar** properties than (coastal & equatorial) **Kelvin waves**
- Dense **essential spectrum** when $0 < |\omega| < \omega$.

INERTIA-GRAVITY WAVES

Elliptic interval (2/2)

Pressure problem

$$\mathcal{P}(\Phi) := -[\omega^2 - N^2]\nabla^2\Phi - N^2(\mathbf{1}_z \cdot \nabla)^2\Phi + (2\boldsymbol{\Omega} \cdot \nabla)^2\Phi = 0$$

→ Principal symbol defines a **Riemannian** metric R on ∂V

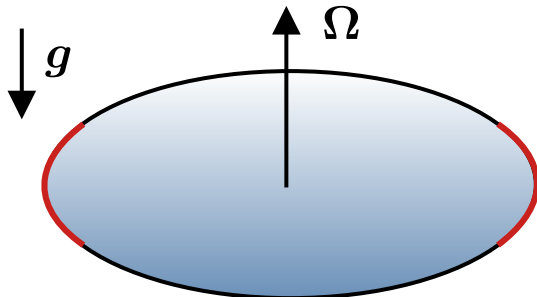
→ Unit normal vector \mathbf{m} w/r R

$$\mathbf{n} \cdot \nabla\Phi|_{\partial V} + \mathcal{B}(\mathbf{n} \times \nabla\Phi) = 0$$

Dirichlet-Neumann operator

$$\begin{aligned} \mathcal{P}(\Phi) = 0 \\ \mathbf{m} \cdot \nabla\Phi|_{\partial V} + \mathcal{V}(\Phi) = 0 \end{aligned} = \boxed{\begin{aligned} \mathcal{N}(\Psi) := \mathbf{m} \cdot \nabla\Phi \\ \mathcal{N}(\Psi) + \mathcal{V}(\Psi) = 0 \end{aligned}} + \begin{aligned} \mathcal{P}(\Phi) = 0 \\ \Phi|_{\partial V} = \Psi \end{aligned}$$

Principal symbol



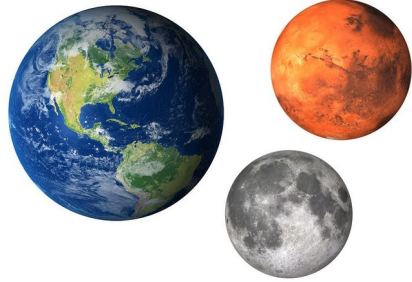
$$\left. \begin{aligned} \sigma(\mathcal{N}) &= \sqrt{\mathfrak{p}} \\ \sigma(\mathcal{V}) &= \mathfrak{v} \end{aligned} \right\} (\sqrt{\mathfrak{p}} + \mathfrak{v})|_{\mathbf{k} \cdot \mathbf{m} = 0} = 0 \quad \text{Shapiro-Lopatinskii condition}$$

Non-elliptical

→ Essential spectrum!

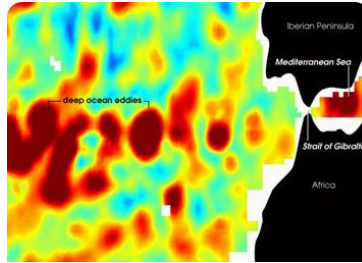
CONCLUSION & PERSPECTIVES

~ 0



Inertial waves

~ 10



Inertia-gravity waves

~ 100



$N/|\Omega|$

Future papers in the series

- **Attractors** of Kelvin waves in other configurations?
- Kelvin waves with **radial gravity**?

Perspectives

